

# Optimal mass transport and density flows

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LCCC, Lund

June 13, 2017

Supported by the NSF & AFOSR

## Plan of the talk:

– Nexus of ideas:

Mass Transport  $\Leftrightarrow$  Schrödinger bridges  $\Leftrightarrow$  Stochastic control  
with a bit on LQG, Riccati, etc.

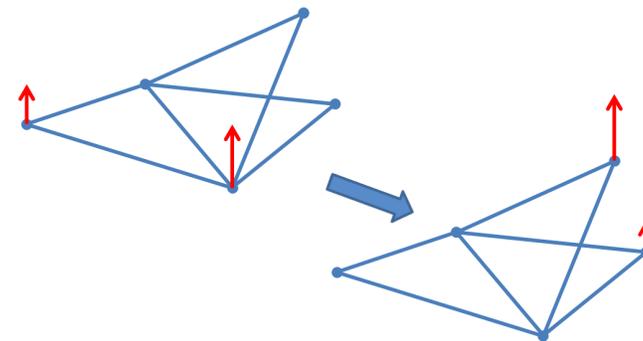
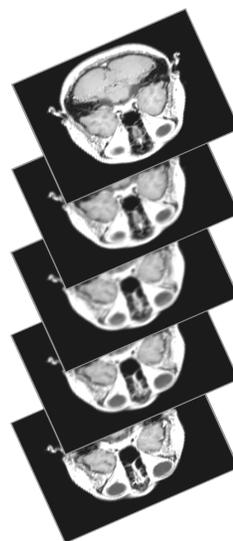
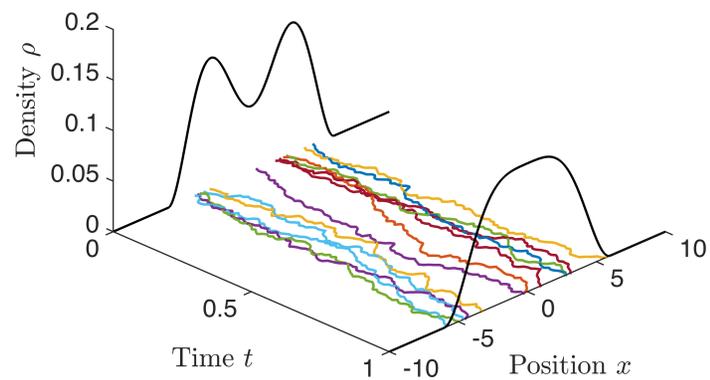
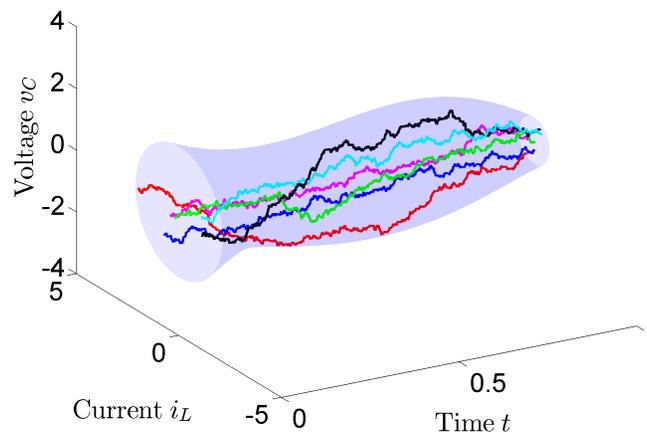
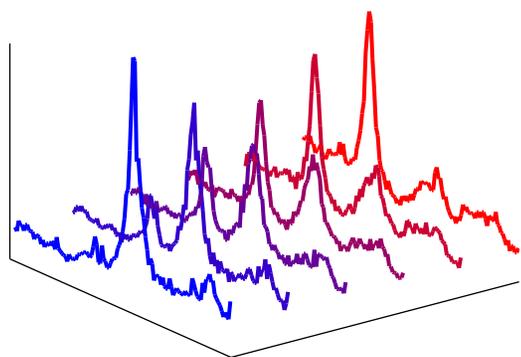
– Discrete-space counterpart:

Markov chains and networks

– Non-commutative counterpart:

Quantum flows & non-commutative geometry

# Density flows



# Optimal Mass Transport (OMT)

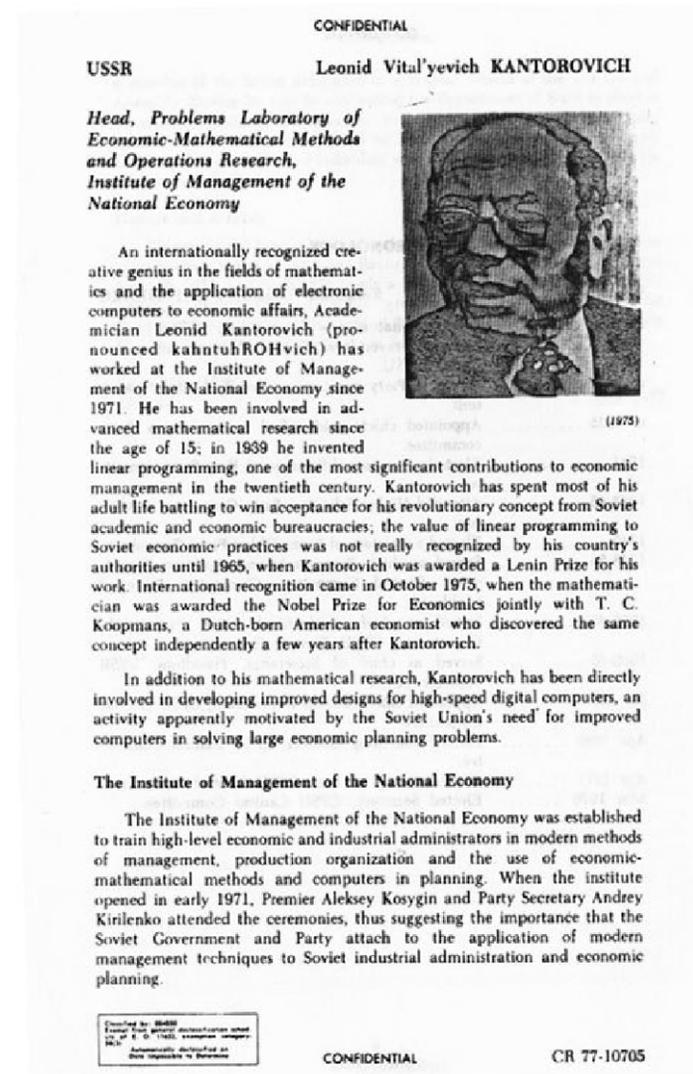


Gaspard Monge 1781



Leonid Kantorovich 1976

Work in early 1940's, Nobel 1975

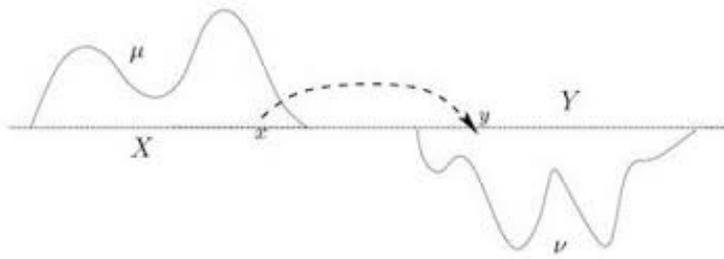


CIA file on Kantorovich  
(wikipedia)

# Monge's formulation

Le mémoire sur les déblais et les remblais

Gaspard Monge 1781



$$\inf_{\mathbf{T}} \int \|x - \underbrace{\mathbf{T}(x)}_y\|^2 d\mu(x)$$



where  $\mathbf{T}\#\mu = \nu$

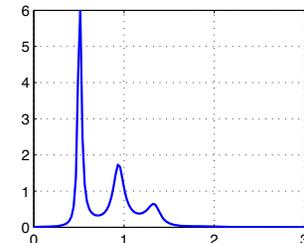
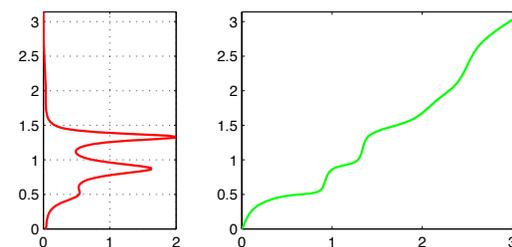
# Kantorovich's formulation

$$\inf_{\pi \in \Pi(\rho_0, \rho_1)} \iint \|x - y\|^2 d\pi(x, y)$$

where  $\Pi(\mu, \nu)$  are “couplings”:

$$\int_y \pi(dx, dy) = \rho_0(x) dx = d\mu(x)$$

$$\int_x \pi(dx, dy) = \rho_1(y) dy = d\nu(y).$$



# B&B's fluid dynamic formulation

Benamou and Brenier (2000):

$$\inf_{(\rho, v)} \int_{\mathbb{R}^n} \int_0^1 \|v(x, t)\|^2 \rho(x, t) dt dx$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (v \rho) = 0$$
$$\rho(x, 0) = \rho_0(x), \quad \rho(y, 1) = \rho_1(y)$$

McCann, Gangbo, Otto, Villani, ...

# Stochastic control formulation

$$\inf_v \mathbb{E}_\rho \left\{ \int_0^1 \|v(x, t)\|^2 dt \right\}$$

$$\dot{x}(t) = v(x, t)$$

$$x(0) \sim \rho_0(x) dx$$

$$x(1) \sim \rho_1(y) dy$$

## OMT as a control problem – derivation

$$\|x - y\|^2 = \inf_{\mathbf{x} \in \mathcal{X}_{xy}} \int_0^1 \|\dot{\mathbf{x}}\|^2 dt,$$

$$\mathcal{X}_{xy} = \{\mathbf{x} \in C^1 \mid \mathbf{x}(0) = x, \mathbf{x}(1) = y\}.$$

Inf attained at constant speed geodesic  $\mathbf{x}^*(t) = (1 - t)x + ty$

## OMT as a control problem – Dirac marginals

Also,  $\text{Inf} =$  any probabilistic average in  $\mathcal{X}_{xy}$

$$\|x - y\|^2 = \inf_{P_{xy}} \mathbb{E}_{P_{xy}} \left\{ \int_0^1 \|\dot{x}(t)\|^2 dt \right\},$$

$P_{xy} \in \mathbb{D}(\delta_x, \delta_y)$  : prob. measures on  $C^1$  with delta marginals

# OMT as a control problem – general marginals

$$\inf_{\pi \in \Pi(\rho_0, \rho_1)} \int_{\mathbb{R}^n \times \mathbb{R}^n} \|x - y\|^2 d\pi(x, y) = \inf_{P \in \mathbb{D}(\rho_0, \rho_1)} \mathbb{E}_P \left\{ \int_0^1 \|\dot{x}(t)\|^2 dt \right\}.$$

⇒ OMT  $\simeq$  stochastic control problem

with atypical boundary constraints

$$\inf_v \mathbb{E} \left\{ \int_0^1 \|v\|^2 dt \right\}$$
$$\dot{x}(t) = v(x(t), t), \quad \text{a.s.}, \quad x(0) \sim \rho_0 dx, \quad x(1) \sim \rho_1 dy.$$

# Schrödinger's Bridges

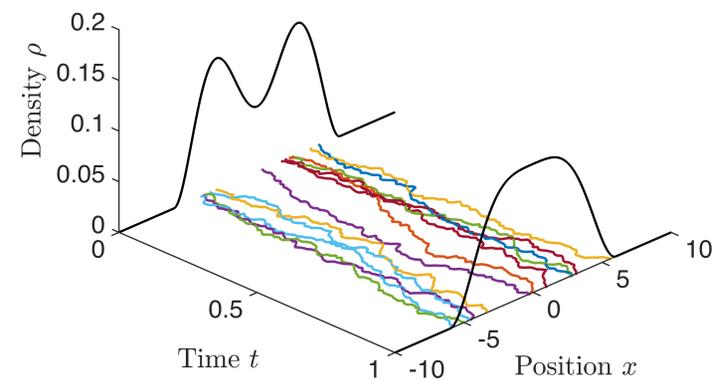


Erwin Schrödinger  
Work in 1926, Nobel 1935

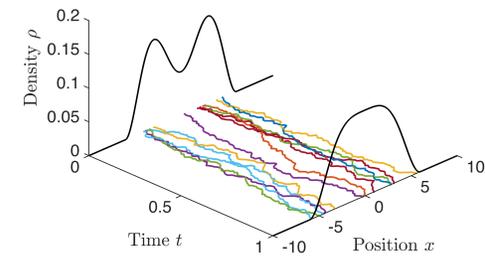
Bridges 1931/32

$$\rho = \Psi \bar{\Psi}$$

$$\Psi_t = U(t)\Psi_0$$



# Schrödinger's Bridge Problem (SBP)



- Cloud of  $N$  independent Brownian particles ( $N$  large)
- empirical distr.  $\rho_0(x)dx$  and  $\rho_1(y)dy$  at  $t = 0$  and  $t = 1$ , resp.
- $\rho_0$  and  $\rho_1$  not compatible with transition mechanism

$$\rho_1(y) \neq \int_0^1 p(t_0, x, t_1, y) \rho_0(x) dx,$$

where

$$p(s, y, t, x) = [2\pi(t - s)]^{-\frac{n}{2}} \exp \left[ -\frac{|x - y|^2}{2(t - s)} \right], \quad s < t$$

Particles have been transported in an unlikely way

**Schrödinger** (1931): Of the many unlikely ways in which this could have happened, which one is the most likely?

# Large deviations formulation of SBP

$$\text{Minimize } H(Q, W) = E_Q \left[ \log \frac{dQ}{dW} \right]$$

over  $Q \in \mathbb{D}(\rho_0, \rho_1)$  distributions on paths with marginals  $\rho$ 's

$H(\cdot, \cdot)$ : relative entropy

**Föllmer 1988**: This is a problem of *large deviations of the empirical distribution* on path space connected through Sanov's theorem to a *maximum entropy problem*.

## Relative entropy w.r.t. Wiener measure

$$dX = vdt + dB$$

Girsanov:

$$E_Q \left[ \log \frac{dQ}{dW} \right] = E_Q \left[ \frac{1}{2} \int_0^t \|v\|^2 ds \right]$$

is a quadratic cost!!!

# SBP as a stochastic control problem

$$\begin{aligned} & \inf_{(\rho, v)} \int_{\mathbb{R}^n} \int_0^1 \|v(x, t)\|^2 \rho(x, t) dt dx, \\ & \frac{\partial \rho}{\partial t} + \nabla \cdot (v \rho) = \frac{1}{2} \Delta \rho \\ & \rho(x, 0) = \rho_0(x), \quad \rho(y, 1) = \rho_1(y). \end{aligned}$$

Blaquière, Dai Pra, ...

compare with OMT:

$$\begin{aligned} & \inf_{(\rho, v)} \int_{\mathbb{R}^n} \int_0^1 \frac{1}{2} \|v(x, t)\|^2 \rho(x, t) dt dx \\ & \frac{\partial \rho}{\partial t} + \nabla \cdot (v \rho) = 0 \\ & \rho(x, 0) = \rho_0(x), \quad \rho(y, 1) = \rho_1(y) \end{aligned}$$

# Fluid-dynamic formulation of SBP

(time-symmetric)

$$\inf_{(\rho, v)} \int_{\mathbb{R}^n} \int_0^1 \left[ \|v(x, t)\|^2 + \left\| \frac{1}{2} \nabla \log \rho(x, t) \right\|^2 \right] \rho(x, t) dt dx,$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (v \rho) = 0,$$
$$\rho(0, x) = \rho_0(x), \quad \rho(1, y) = \rho_1(y).$$

$\left\| \frac{1}{2} \nabla \log \rho(x, t) \right\|^2$  : Fisher information, Nelson's osmotic power

Chen-Georgiou-Pavon, On the relation between optimal transport and Schrödinger bridges: A stochastic control viewpoint, *J. Opt. Theory Appl.*, 2015

Mikami 2004, Mikami-Thieullen 2006,2008, Léonard 2012

# Erwin Schrödinger's insight on SBP

the density factors into

$$\rho(x, t) = \varphi(x, t)\hat{\varphi}(x, t)$$

where  $\varphi$  and  $\hat{\varphi}$  solve (Schrödinger's system):

$$\begin{aligned}\varphi(x, t) &= \int p(t, x, 1, y)\varphi(y, 1)dy, & \varphi(x, 0)\hat{\varphi}(x, 0) &= \rho_0(x) \\ \hat{\varphi}(x, t) &= \int p(0, y, t, x)\hat{\varphi}(y, 0)dy, & \varphi(x, 1)\hat{\varphi}(x, 1) &= \rho_1(x).\end{aligned}$$

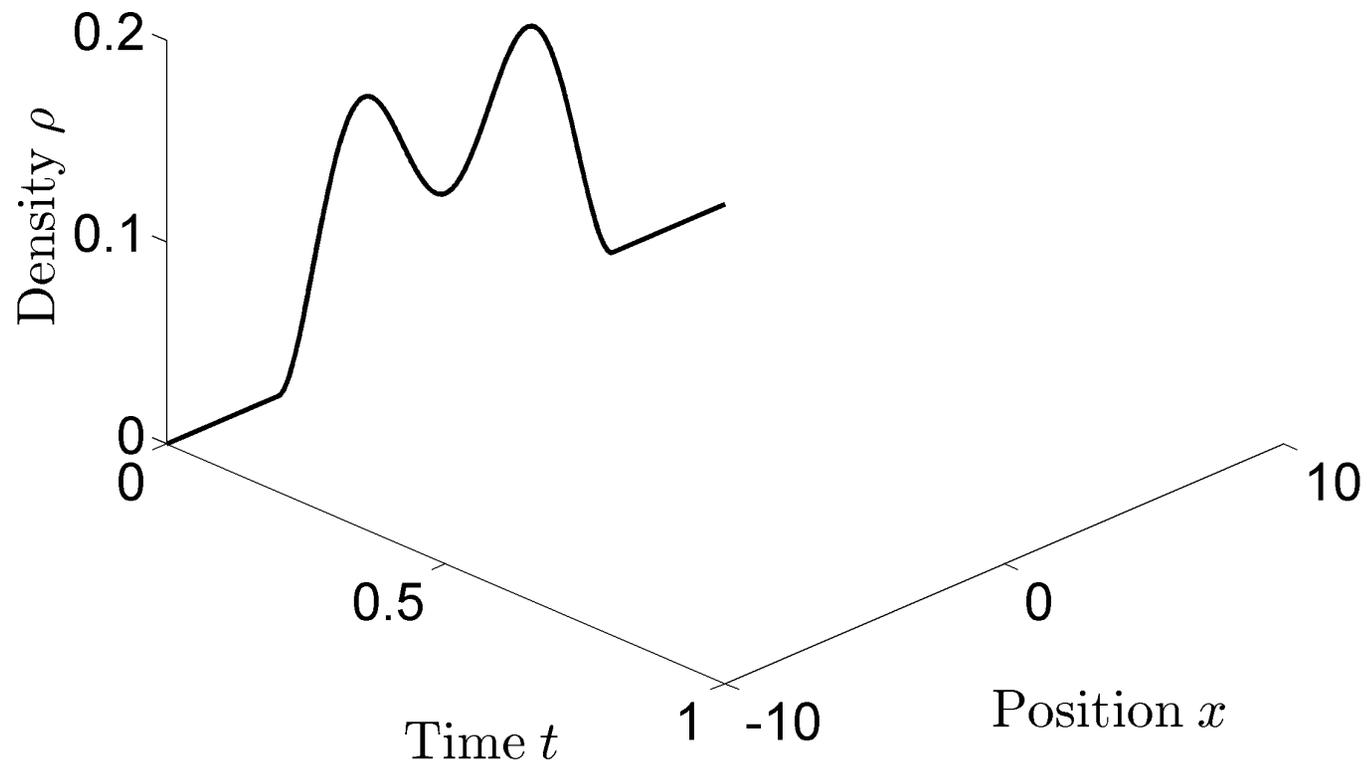
compare with  $\Psi\bar{\Psi} = \rho$

Existence and uniqueness for Schrödinger's system:

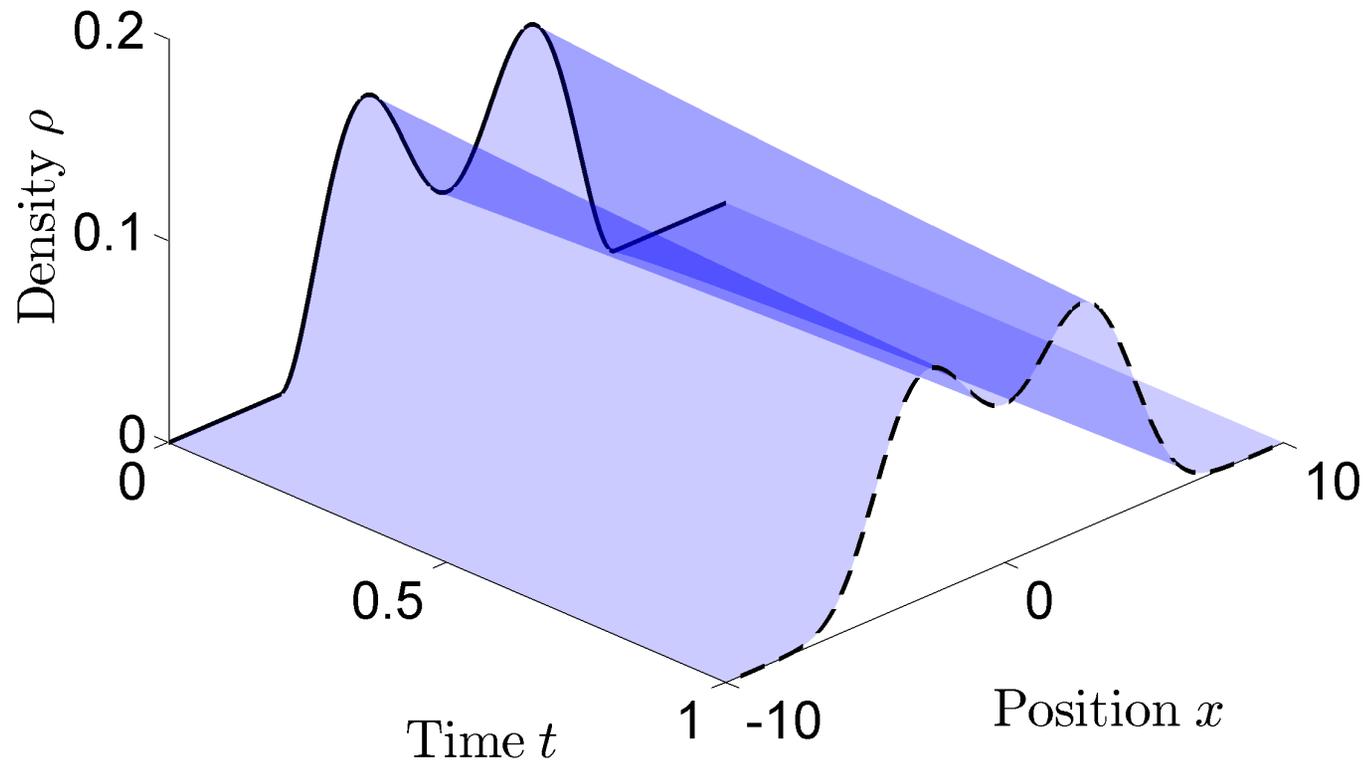
Fortet 1940, Beurling 1960, Jamison 1974/75, Föllmer 1988.

~ Sinkhorn iteration & Quantum version: Georgiou-Pavon 2015

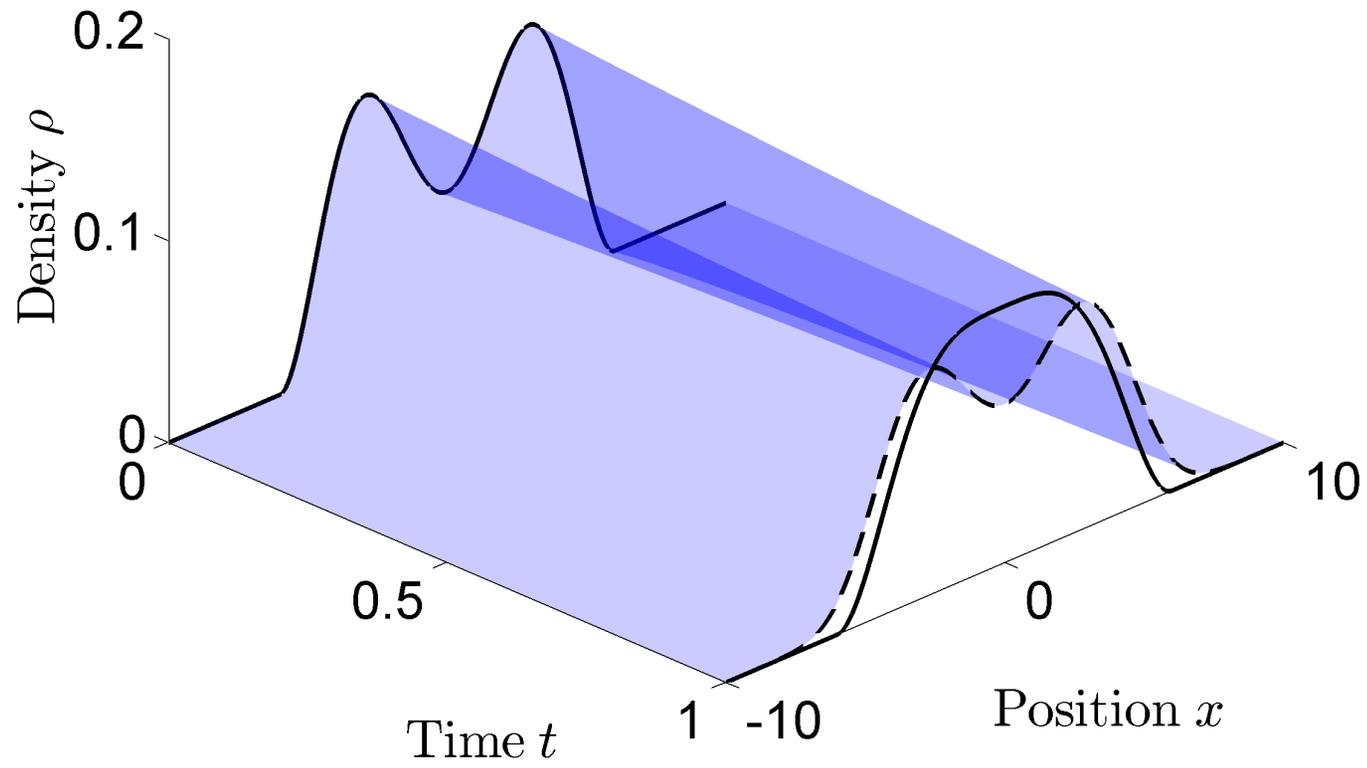
# SBP schematic



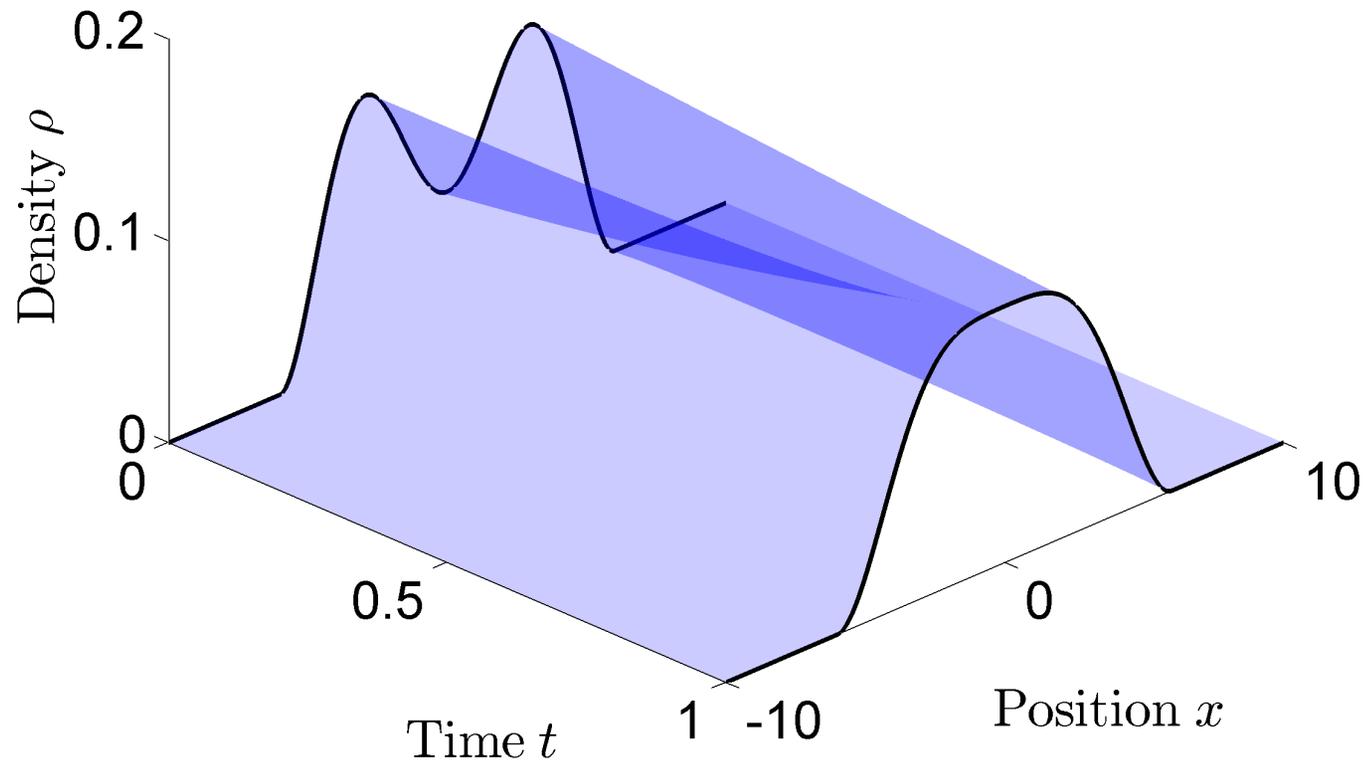
# SBP schematic



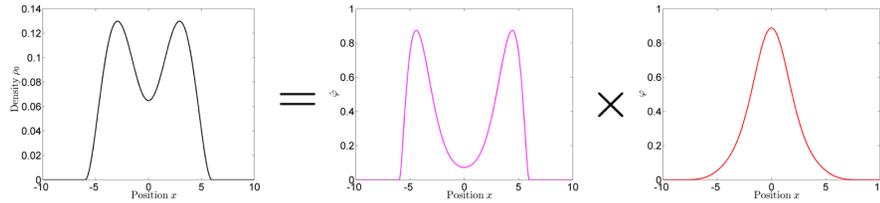
# SBP schematic



# SBP schematic

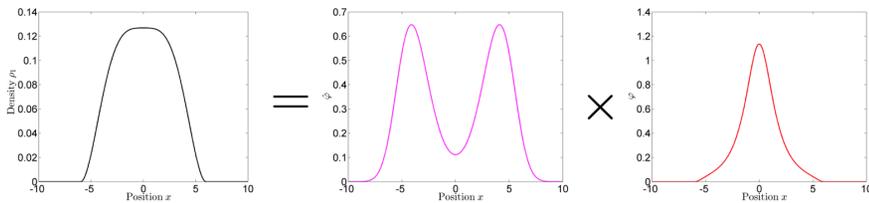


# Schrödinger system



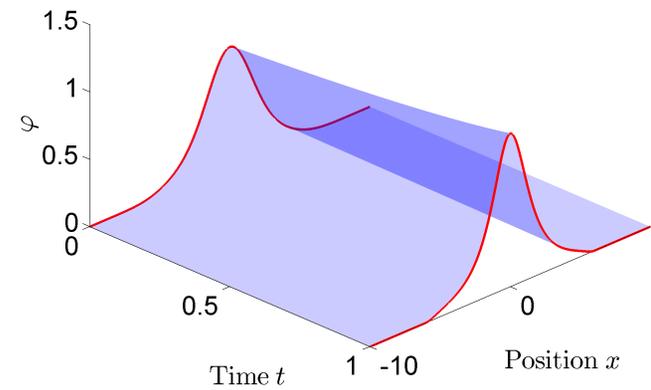
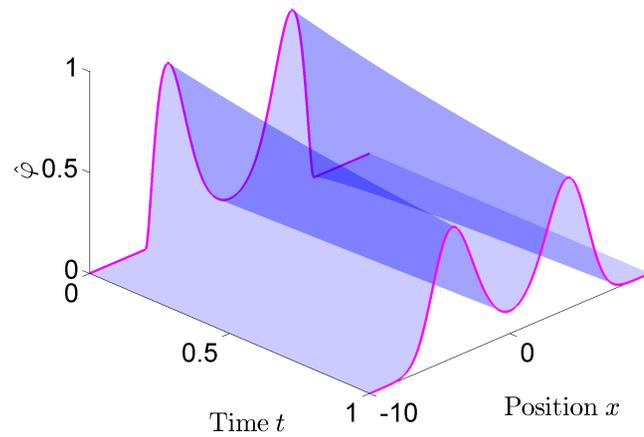
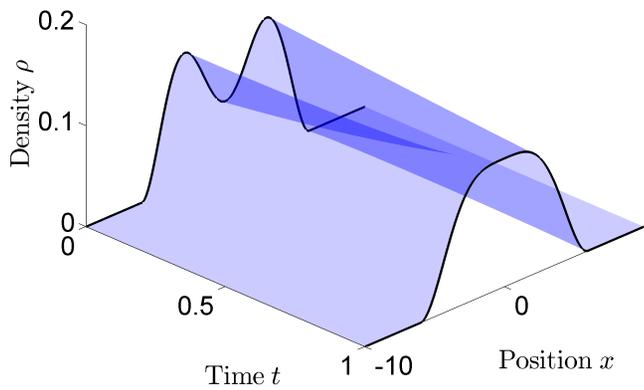
$$-\frac{\partial \varphi}{\partial t}(t, x) = \frac{1}{2} \Delta \varphi(t, x)$$

$$\frac{\partial \hat{\varphi}}{\partial t}(t, x) = \frac{1}{2} \Delta \hat{\varphi}(t, x)$$



$$\varphi(0, x) \hat{\varphi}(0, x) = \rho_0(x)$$

$$\varphi(1, x) \hat{\varphi}(1, x) = \rho_1(x)$$



# Existence & uniqueness (Sinkhorn scaling)

$$\begin{array}{ccc}
 \begin{array}{c} \hat{\varphi} \\ (\rho_0) \uparrow \\ \varphi \end{array} & \begin{array}{c} \xrightarrow{\Delta/2} \\ \\ \xleftarrow{-\Delta/2} \end{array} & \begin{array}{c} \hat{\varphi} \\ \downarrow (\rho_1) \\ \varphi \end{array}
 \end{array}
 \quad
 \begin{array}{l}
 -\frac{\partial \varphi}{\partial t}(t, x) = \frac{1}{2} \Delta \varphi(t, x) \\
 \frac{\partial \hat{\varphi}}{\partial t}(t, x) = \frac{1}{2} \Delta \hat{\varphi}(t, x) \\
 \varphi(0, x) \hat{\varphi}(0, x) = \rho_0(x) \\
 \varphi(1, x) \hat{\varphi}(1, x) = \rho_1(x)
 \end{array}$$

iteration is **contractive** in the Hilbert metric!

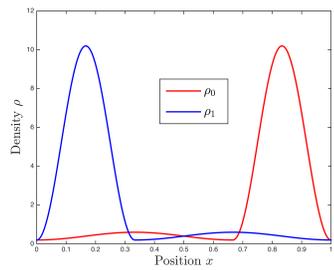
$$d_H(p, q) = \log \frac{M(p, q)}{m(p, q)}$$

$$M(p, q) := \inf \{ \lambda \mid p \leq \lambda q \}$$

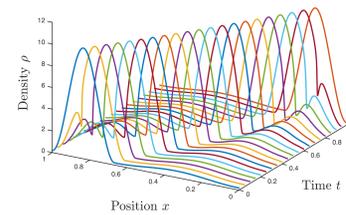
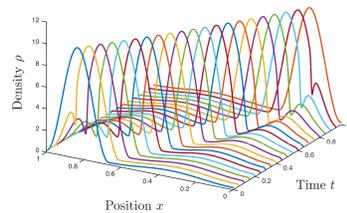
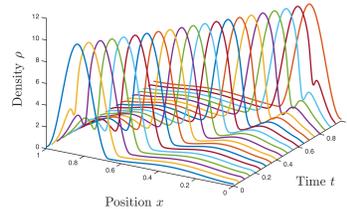
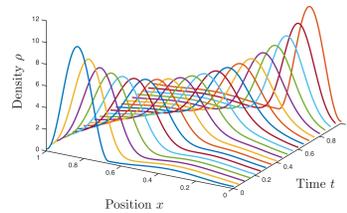
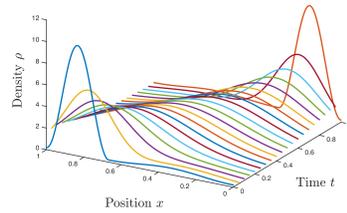
$$m(p, q) := \sup \{ \lambda \mid \lambda q \leq p \}$$

Chen-Georgiou-Pavon, Entropic and displacement interpolation: a computational approach using the Hilbert metric, *SIAM J. Appl. Math.*

# OMT as limit to SBP: numerics in general



Marginal distributions



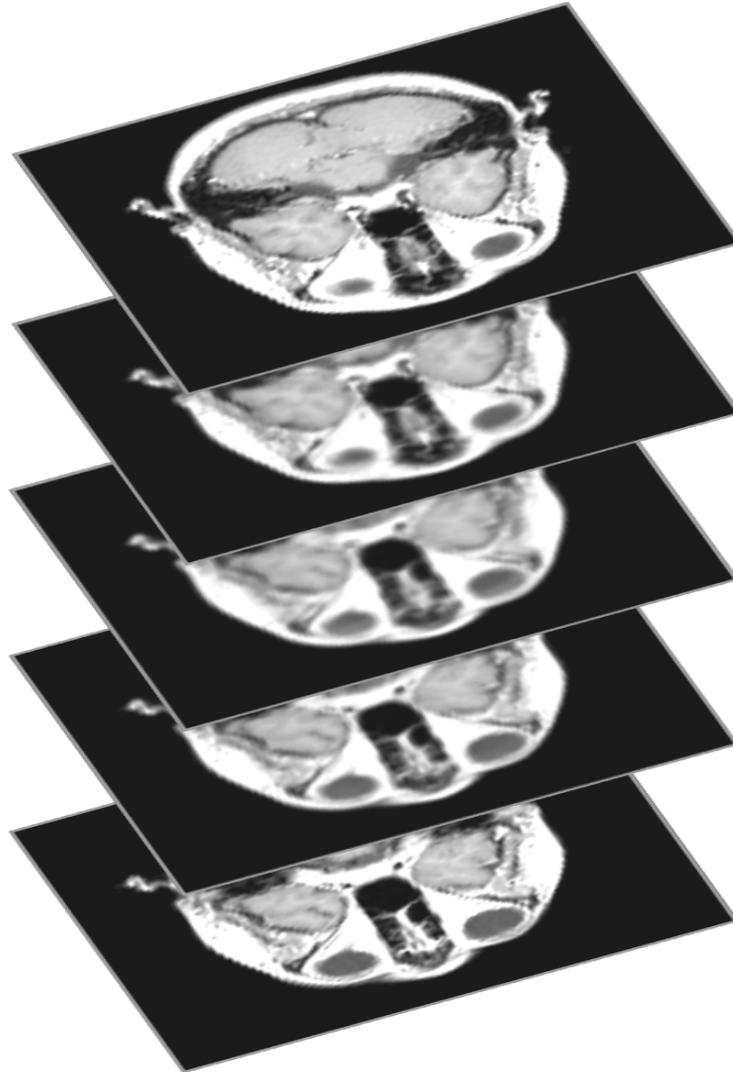
OMT interpolation:

$$\rho_t + \nabla \cdot \rho v = 0$$

$$\rho_t + \nabla \cdot \rho v = \epsilon \Delta \rho, \text{ varying } \epsilon$$

# Applications: Image interpolation

Interpolation of 2D images to a 3D model:



# LQG - covariance control

$$\min_u \mathbb{E} \left\{ \int_0^T \|u(t)\|^2 dt \right\},$$

s.t.

$$dX = AXdt + Budt + BdW$$

$$X(0) \sim \mathcal{N}(0, \Sigma_0), \quad X(T) \sim \mathcal{N}(0, \Sigma_1) \quad \leftarrow \text{these are the } \rho\text{'s}$$

Beghi (1996), Grigoriadis- Skelton (1997)

Brockett (2007, 2012), Vladimirov-Petersen (2010, 2015)

# Bridges - LQG - covariance control in general

$$\min_u \mathbb{E} \left\{ \int_0^T \|u(t)\|^2 dt \right\},$$

s.t.

$$dX = AXdt + Buds + B_1dW$$

$$X(0) \sim \mathcal{N}(0, \Sigma_0), \quad X(T) \sim \mathcal{N}(0, \Sigma_1)$$

connection with SBP  $\Rightarrow \phi(t, x) = \exp(-\|x\|_{Q(t)}^2)$  & Riccati's

# SBP Riccati's

– nonlinearly coupled Riccati equations  $\equiv$  Schrödinger system

$$\dot{\Pi} = -A'\Pi - \Pi A + \Pi B B' \Pi$$

$$\begin{aligned} \dot{\mathbf{H}} = & -A'\mathbf{H} - \mathbf{H}A - \mathbf{H}B B' \mathbf{H} \\ & + (\Pi + \mathbf{H}) (B B' - B_1 B_1') (\Pi + \mathbf{H}). \end{aligned}$$

$$\Sigma_0^{-1} = \Pi(0) + \mathbf{H}(0)$$

$$\Sigma_T^{-1} = \Pi(T) + \mathbf{H}(T).$$

$$\log(\rho) = \log(\phi) + \log(\hat{\phi}) \Leftrightarrow \Sigma^{-1} = \Pi + \mathbf{H}$$

Chen-Georgiou-Pavon, Optimal steering of a linear stochastic system to a final probability distribution, *IEEE Trans. Aut. Control*, May 2016

## stationary SBP

When can  $\Sigma$  be a stationary state-covariance for

$$dx(t) = (A - BK)x(t)dt + B_1dw(t)?$$

i.e., when is  $\Sigma = Exx'$ , for suitable choice of  $K$ ?

– not all  $\Sigma$  can be realized by state feedback

# stationary SBP

When can  $\Sigma$  be a stationary state-covariance for

$$dx(t) = (A - BK)x(t)dt + B_1dw(t)?$$

This is so iff

$$\text{rank} \begin{bmatrix} A\Sigma + \Sigma A' + B_1B_1' & B \\ B & 0 \end{bmatrix} = \text{rank} \begin{bmatrix} 0 & B \\ B & 0 \end{bmatrix}.$$

- Chen-Georgiou-Pavon, Optimal steering..., Part II *IEEE TAC*, May 2016
- Georgiou, Structure of state covariances... *TAC* 2002
- recent work with Mihailo Jovanovic et al. on inverse problems, etc., 2016, 2017

# stationary SBP

Assuming

$$\text{rank} \begin{bmatrix} A\Sigma + \Sigma A' + B_1 B_1' & B \\ B & 0 \end{bmatrix} = \text{rank} \begin{bmatrix} 0 & B \\ B & 0 \end{bmatrix},$$

find  $K$  so that

for  $u = -Kx$  and  $dx = (A - BK)xdt + B_1 dw$ ,

we have:

$$\Sigma = Exx' \text{ and } J_{\text{power}}(u) := \mathbb{E}\{\|u\|^2\} \text{ is minimal}$$

Via semidefinite programming:

– Chen-Georgiou-Pavon, Optimal steering..., Part II *IEEE TAC*, May 2016.

# Application: Cooling

Efficient steering from initial condition  $\rho_0$  to  $\rho_1$  at finite time

- Efficient stationary state of stochastic oscillators to **desired**  $\rho_1$
- thermodynamic systems, controlling collective response
- magnetization distribution in NMR spectroscopy,..

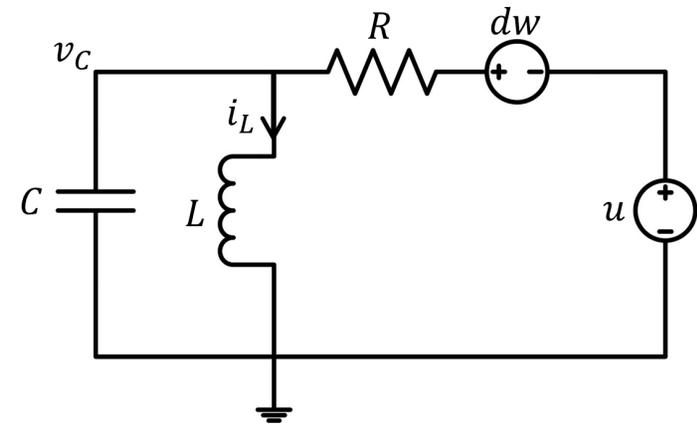
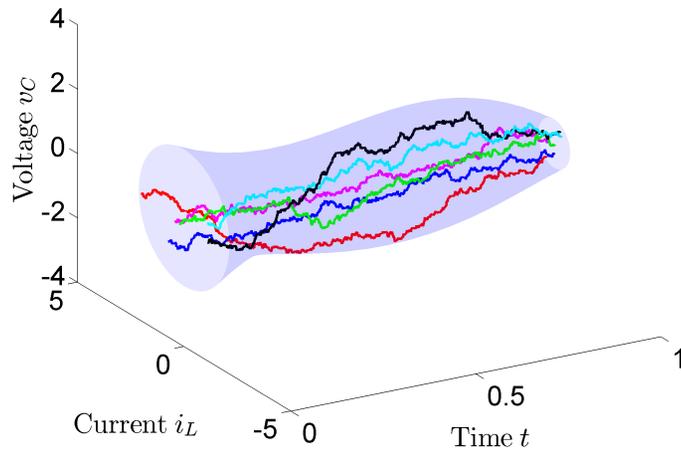
- Chen-Georgiou-Pavon

Fast cooling for a system of stochastic oscillators, *J. Math. Phys.* Nov. 2015.

# Cooling (cont'd)

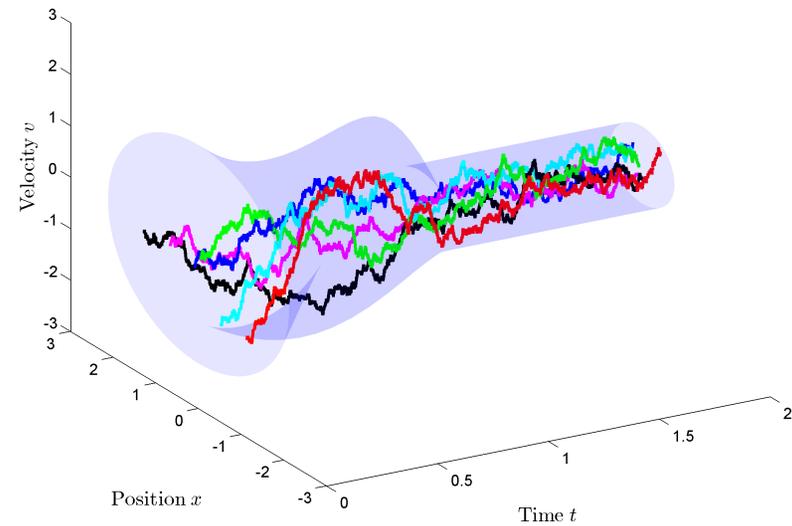
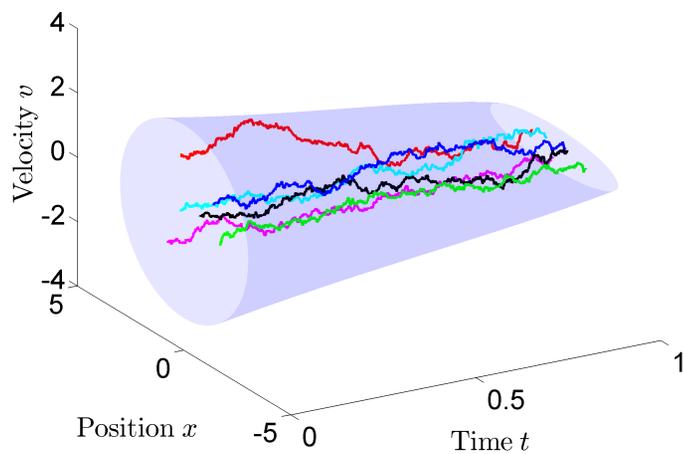
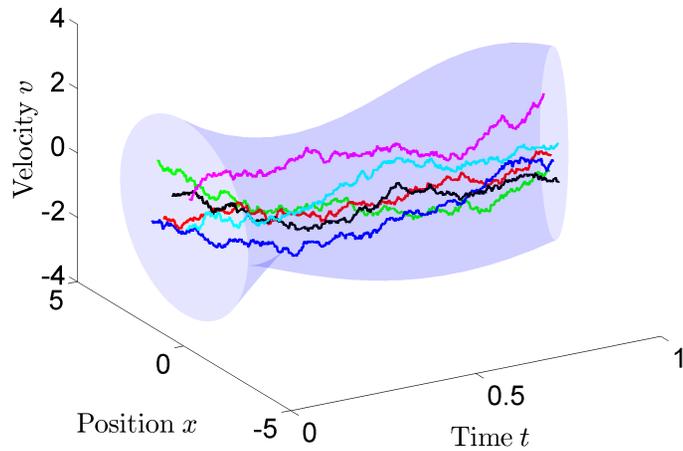
Nyquist-Johnson noise driven oscillator

$$L di_L(t) = v_C(t) dt$$
$$RC dv_C(t) = -v_C(t) dt - Ri_L(t) dt + u(t) dt + dw(t)$$



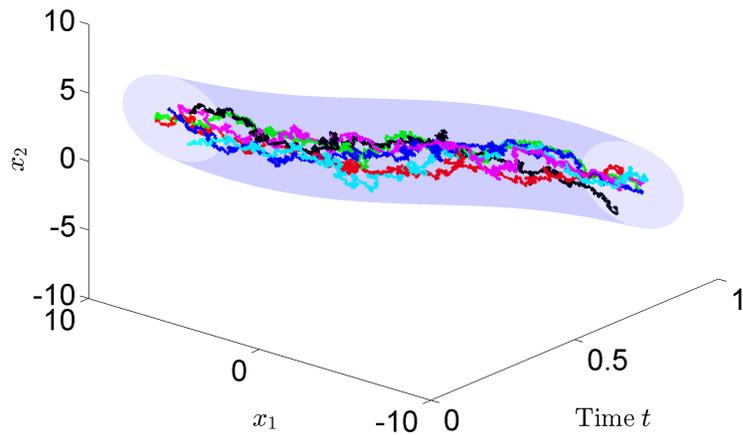
# Cooling & keeping it cool!

Inertial particles with stochastic excitation

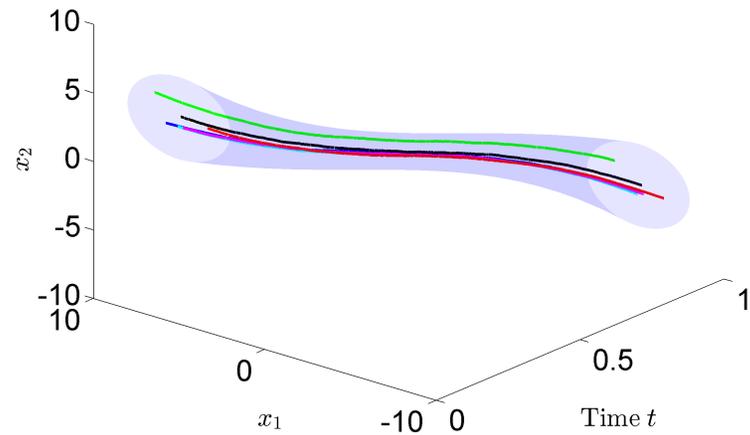


trajectories in phase space  
transparent tube: “ $3\sigma$  region”

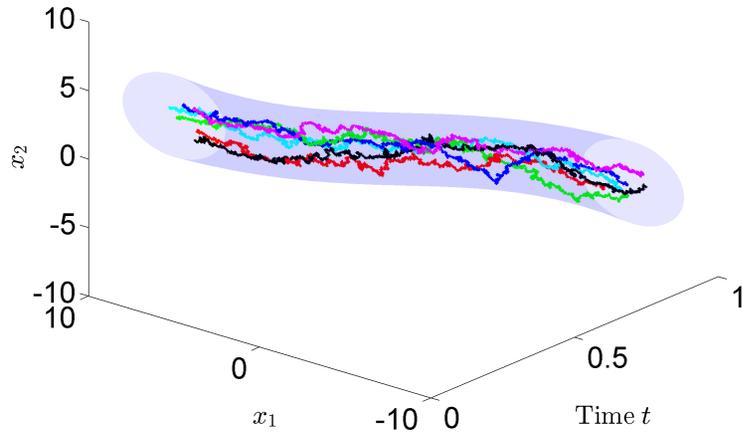
# Application: OMT with dynamics via SBP



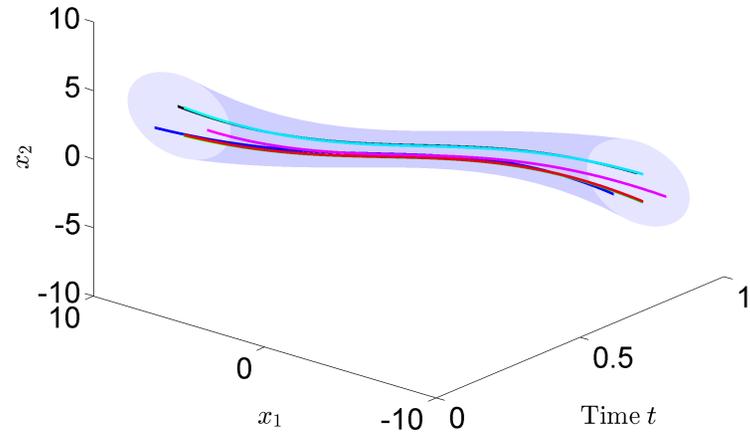
Schrödinger bridge with  $\epsilon = 9$



Schrödinger bridge with  $\epsilon = 0.01$

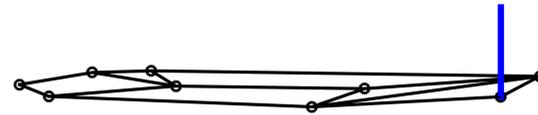
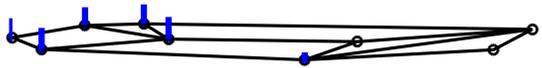
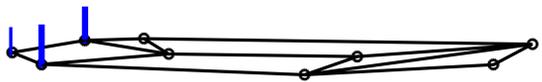
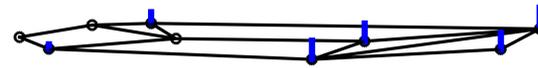
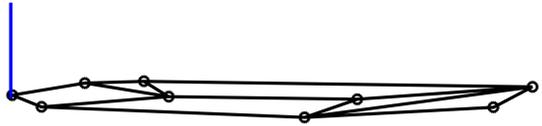


Schrödinger bridge with  $\epsilon = 4$



Optimal transport with prior

# Discrete space: SBP and OMT on graph



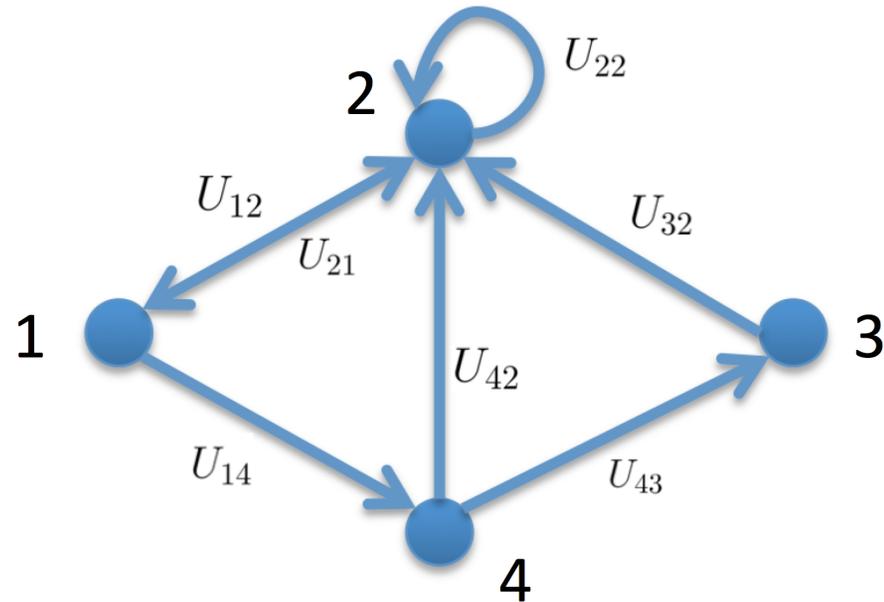
Chen-G-Pavon-Tannenbaum, Robust transport over networks  
TAC to appear in March

# Flow on Graphs - transportation

## Graph:

nodes  $\mathcal{X} = \{1, 2, 3, 4\}$   
edges  $\mathcal{E} = \{(1, 2), (1, 4), \dots\}$

paths:  $(1, 4, 2), (1, 4, 3, 2), \dots$   
transport cost  $i \rightarrow j$ :  $U_{ij}$



## Markov chain

transition mechanism  $Q_{ij}$   $\text{Prob}(i \rightarrow j)$   
portion of mass passing from  $i$  to  $j$ .

$\rho(t, x)$   $\text{Prob}(X(t) = x)$   
“mass” at time  $t$  sitting at node  $x$

$\rho_0(x_0)Q_{x_0x_1} \cdots Q_{x_{N-1}x_N}$   $\text{Prob}(X(0) = x_0 X(1) = x_1 \cdots X(N) = x_N)$   
portion of mass traveling along  $x_0 x_1 \cdots x_N$

# Schrödinger bridges on graphs

initial & final distributions:  $\rho_0$  and  $\rho_N$

transition probabilities  $Q_{ij}$  (prior)

that may not be consistent with the marginals, i.e.,

$$\rho_N(x_N) \neq \sum_{x_0, x_1, \dots, x_{N-1}} \rho_0(x_0) Q_{x_0 x_1} Q_{x_1 x_2} \cdots Q_{x_{N-1} x_N}$$

Determine

$$P^{\text{opt}} = \operatorname{argmin}\{H(P, Q) \mid P \in \mathbb{D}(\rho_0, \rho_N)\}$$

$H(P, Q)$  relative entropy/KL divergence

- Choice of prior influences transport properties!  
Ruelle-Bowen transition probabilities  $\Rightarrow$  “equalize” usage of all alternative paths  
 $\Rightarrow$  dispersive transportation, robustness
- Computations via iterative scaling (Sinkhorn-like)

# Trading cost vs. robustness

cost  $U_{ij}$  in traversing edge  $(i, j)$ :

$$U(x_0, x_1, \dots, x_N) = \sum_{t=0}^{N-1} U_{x_t x_{t+1}}$$

cost of transportation:

$$\mathcal{U}(P) := \sum_{\{(x_0 \dots x_N)\}} P(x_0, x_1, \dots, x_N) U(x_0, x_1, \dots, x_N)$$

minimize “free energy”

$$\mathcal{F}(P) := \mathcal{U}(P) - T\mathcal{S}(P)$$

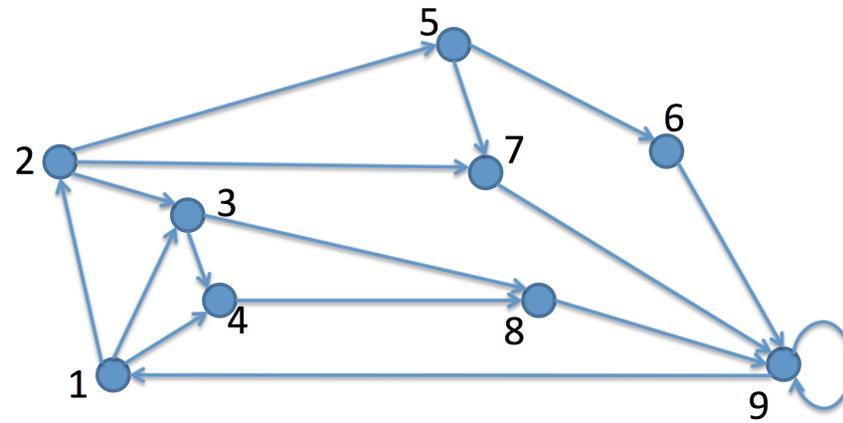
$$= - \sum_{\{\text{paths}\}} P \log(e^{-U}) + T \sum P \log(P) = T H(P | e^{-U/T})$$

“temperature”  $T$ : tradeoff between  $\mathcal{U}$  and  $\mathcal{S}$

i.e., tradeoff between cost & “dispersiveness/robustness”

# Application: transport scheduling

Move a unit mass  
from node **1** to node **9**  
in three steps:



Available paths

$$\left\{ \begin{array}{l} (1 - 2 - 7 - 9) \\ (1 - 3 - 8 - 9) \\ (1 - 4 - 8 - 9) \end{array} \right.$$

Solution  $\rho(t, x)$ : equal usage of the three options.

# Matrix-valued OMT & SBP

## our goal

extend the fluid dynamics framework to

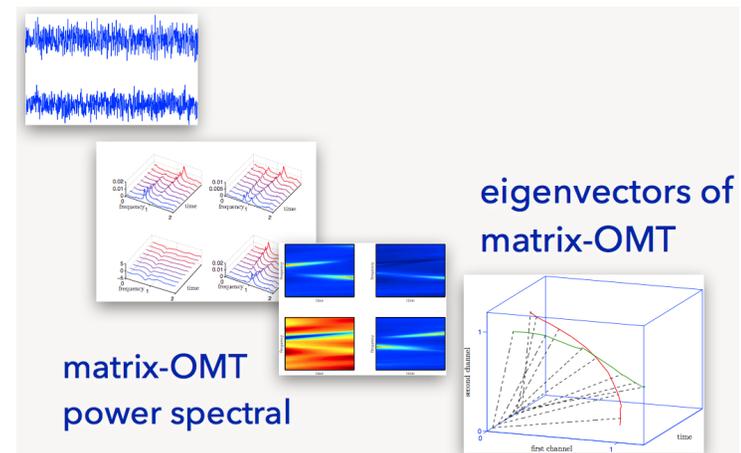
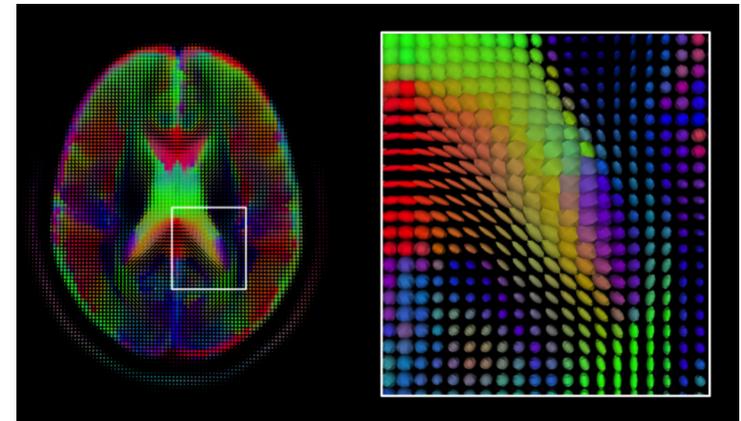
- Hermitian matrices
- matrix-valued distributions

I.e., formulate for matrices...

$$\inf \int_{\text{space}} \int_0^1 \rho(t, x) \|v(t, x)\|^2 dt dx$$

$$\frac{\partial \rho}{\partial t} + \nabla_x \cdot (\rho v) = 0$$

$$\rho(0, \cdot) = \rho_0, \rho(1, \cdot) = \rho_1$$



# Quantum mechanics

Starting point: Lindblad equation

Hermitian  $\geq 0$  (trace 1): density matrices

$$\begin{aligned}\dot{\rho} = & -[iH, \rho] \\ & + \sum_{k=1}^N (L_k \rho L_k - \frac{1}{2} \rho L_k L_k - \frac{1}{2} L_k L_k \rho),\end{aligned}$$

compare with

$$\rho_t = -\nabla \cdot (\rho v)$$

# Some calculus

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for ordinary functions:

$$f(x) : g(x) \mapsto f(x)g(x)$$

$$\partial_x : g(x) \mapsto \partial_x g(x)$$

$$[\partial_x, f(x)] : g(x) \mapsto \partial_x f(x)g(x) - f(x)\partial_x g(x) = (\partial_x f(x))g(x)$$

---

For matrices:

$$\partial_{L_i} X = [L_i, X] = [L_i X - X L_i] \quad \text{and} \quad \nabla_L : X \mapsto \begin{bmatrix} L_1 X - X L_1 \\ \vdots \\ L_N X - X L_N \end{bmatrix}$$

and more calculus!

$\nabla_L$  satisfies

$$\nabla_L(XY) = (\nabla_L X)Y + X(\nabla_L Y)$$

*divergence:*

$$\nabla_L^* : \mathcal{S}^N \rightarrow \mathcal{H}, Y = \begin{bmatrix} Y_1 \\ \vdots \\ Y_N \end{bmatrix} \mapsto \sum_k^N L_k Y_k - Y_k L_k.$$

i.e., using  $\langle X, Y \rangle = \sum_{k=1}^N \text{tr}(X_k^* Y_k)$

$$\langle \nabla_L X, Y \rangle = \langle X, \nabla_L^* Y \rangle$$

Back to Lindblad's equation!

$$\dot{\rho} + \nabla_{iH}^* \rho = \underbrace{-\nabla_L^* \nabla_L \rho}_{-\Delta \rho}$$

Schrödinger's equation:  $\dot{\rho} + \nabla_{iH}^* \rho = 0$

# Matrix continuity equation in general

$$\dot{\rho} + \nabla_L^*(\rho \circ v) = 0$$

choices of non-commutative “momentum”

$$(\rho \circ v) =$$

$$\frac{1}{2}(\rho v + v \rho) \quad (\text{“anti-commutator”})$$

$$\int_0^1 \rho^s v \rho^{1-s} ds \quad (\text{Kubo-Mori})$$

$$\rho^{1/2} v \rho^{1/2}$$

# Matrix OMT (... and SB)

$$\min_{\rho, v} \int_0^1 \text{tr}(\rho v^* v) dt,$$

$$\dot{\rho} = \frac{1}{2} \nabla_L^* (\rho v + v \rho),$$

$$\rho(0) = \rho_0, \quad \rho(1) = \rho_1,$$

$v$ : vector of matrices

$$v^* v = \sum_{k=1}^N v_k^* v_k$$

$L$ : “non-commutative coordinates” (in place of  $x_1, x_2, \dots$ )

# matrix-OMT geometry, gradient flows, and more

Quantum:

Lindblad's equation = gradient flow of the von Neumann entropy

Yongxin Chen, TTG & Allen Tannenbaum

“Matrix OMT: a Quantum Mechanical approach,” 2016

Eric Carlen & Jan Maas

“Gradient flow and entropic inequalities...,” 2016

Markus Mittnenzweig & Alexander Mielke

“An entropic gradient structure for Lindblad...,” 2016.

Yongxin Chen, W. Gangbo, TTG & Allen Tannenbaum

“On the matrix Monge-Kantorovich”

# Gradient flow of Entropy

$$\begin{aligned}\frac{dS(\rho(t))}{dt} &= \dots \\ &= -\text{tr}((\nabla_L \log \rho)^* \int_0^1 \rho^s v \rho^{1-s} ds),\end{aligned}$$

⇒ greatest ascent direction  $v = -\nabla_L \log \rho$ .

---

non-commutative analog of:  $\partial_x \rho = \rho \partial_x (\log \rho)$ :

$$\nabla_L \rho = \int_0^1 \rho^s (\nabla_L \log \rho) \rho^{1-s} ds$$

---

**Gradient flow:**

$$\dot{\rho} = -\nabla_L^* \int_0^1 \rho^s (\nabla_L \log \rho) \rho^{1-s} ds = -\nabla_L^* \nabla_L \rho = \Delta_L \rho,$$

Linear heat equation (now Lindblad) just as in the scalar case!

# matrix-OMT geometry, gradient flows, and more

Quantum:

Lindblad's equation = gradient flow of the von Neumann entropy

Medical imaging (DTI imaging):

matricial geodesics

Time series (matrix-spectrograms):

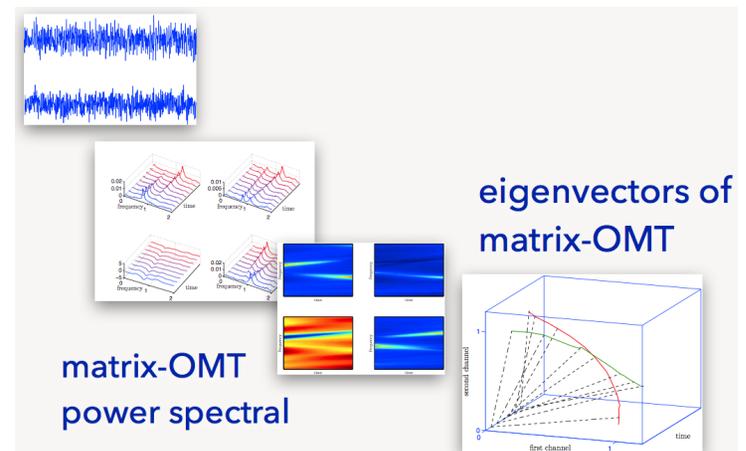
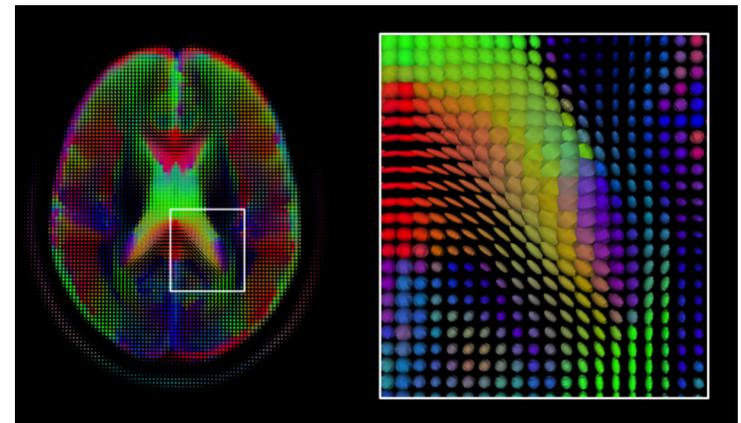
non-stationary processes

arXiv 2016:

Chen-G-Tannenbaum

Chen-Gangbo-G-Tannenbaum

also Carlen-Maas, Mittnenzweig-Mielke



# Concluding remarks

Control problem:

steering flow between specified marginals

Modeling/interpolation problem:

reconciling flow with the known prior

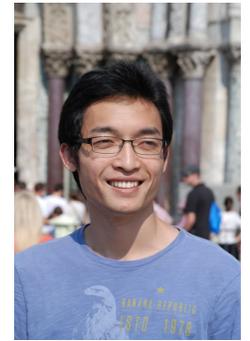
Metrics, metrics, metrics:

for interpolation, smoothing, etc.

## Applications:

- time-series analysis, spectral flows,...(original motivation)
- control of collective motion of particles, agents,..
- transportation of resources with end-point specs
- tradeoffs between cost and robustness in transport problems
- thermodynamics, quantum

Thank you for your attention



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Michele Pavon



Wilfrid Gangbo



Allen Tannenbaum